#### Emergent topology of insulators Noncommutative index theory Warsaw

Terry A. Loring

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Where string hits D-brane is not precisely determined. Shown: fuzzy sphere, fuzzy torus.

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It's square is almost the Laplace operator:

$$\mathcal{D}_{\mathbf{0}}^{2} = \begin{bmatrix} X^{2} + Y^{2} + Z^{2} & 0\\ 0 & X^{2} + Y^{2} + Z^{2} \end{bmatrix} + \begin{bmatrix} i [Y, Z] & [X, Y] - i [X, Z]\\ - [X, Y] - i [X, Z] & -i [Y, Z] \end{bmatrix}$$

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So the emergent surface is what mathematicians call the Joint Clifford Spectrum.

#### Definition

For Hermitian operators X, Y, Z the Clifford spectrum is the set

$$\Lambda(X,Y,Z) = \left\{ \lambda \in \mathbb{R}^d \, | \, D_\lambda(X,Y,Z) \text{ is singular} \right\}.$$

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$$\Delta_{\lambda}(X,Y,Z) = (X - \lambda_1)^2 + (Y - \lambda_2)^2 + (Z - \lambda_3)^2$$

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In the case of *commuting* hermitian X, Y and Z,

 $D_{\lambda}(X, Y, Z)$  is singular  $\iff \Delta_{\lambda}(X, Y, Z)$  is singular  $\iff \exists \mathbf{v} \text{ a unit vector with } X\mathbf{v} = \lambda_1 \mathbf{v}, \ Y\mathbf{v} = \lambda_2 \mathbf{v}, \ Z\mathbf{v} = \lambda_3 \mathbf{v}$ 

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This means  $\Lambda(X, Y, Z)$  is a finite set.

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$$X_{n} = \frac{1}{n} \begin{bmatrix} n & & & \\ & n-1 & & \\ & & \ddots & \\ & & & 1-n & \\ & & & & -n \end{bmatrix},$$

$$Y_n = \frac{1}{2n} (T_n^* + T_n), \quad Z_n = \frac{i}{2n} (T_n^* - T_n)$$

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where we set  $N_n = n(n+1)$  and

$$T_n = \begin{bmatrix} 0 & \sqrt{N_n - (n-1)n} & & & \\ & 0 & \sqrt{N_n - (n-2)(n-1)} & & \\ & & \ddots & \ddots & \\ & & 0 & \sqrt{N_n - (-n)(1-n)} \\ & & & 0 \end{bmatrix}$$

The second example is  $X_n$ ,  $\tilde{Y}_n$ ,  $\tilde{Z}_n$  where

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$$S_n = \left[ egin{array}{ccccc} 0 & 1 & & & \ & 0 & 1 & & \ & & \ddots & \ddots & \ & & & 0 & 1 \ & & & & 0 \end{array} 
ight].$$

We'll see  $\Lambda(X, Y, Z)$  and  $\Lambda(X, \tilde{Y}, \tilde{Z})$  are uncountable sets.

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#### Definition

For Hermitian operators *X*, *Y*, *Z* and  $\epsilon \ge 0$ , define the *Clifford*  $\epsilon$ -pseudospectrum to be the set

$$\Lambda_{\epsilon}(X,Y,Z) = \left\{ \lambda \in \mathbb{R}^d \, \middle| \, \left\| \mathcal{D}_{\lambda}(X,Y,Z)^{-1} \right\|^{-1} \leq \epsilon \right\}$$

with the convention that *R* singular means  $||R^{-1}||^{-1} = 0$ .

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Notice

$$\left\| \mathcal{D}_{\lambda}(X,Y,Z)^{-1} \right\|^{-1} = \left| \operatorname{eig}_{\min} \left( \mathcal{D}_{\lambda}(X,Y,Z) \right) \right|$$

called the gap at  $\lambda$ .



**Figure :** Left: approximation of part of the Clifford spectrum for the usual fuzzy sphere, spin=2. Right: Same spin but with modified Y and Z.



Figure : Now spin=20.



#### Figure : Now spin=200.



#### Figure : Now spin=2000.

# Obviously $\Lambda(X, Y, Z)$ is a subset of $\Lambda_{\epsilon}(X, Y, Z)$ . Could be a finite set?
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#### Definition

The *local index* for the finite system (X, Y, Z) is then

$$\operatorname{Ind}_{\lambda}(X, Y, Z) = \frac{1}{2}\operatorname{Sig}(\mathcal{D}_{\lambda}(X, Y, Z))$$

defined at every point where  $D_{\lambda}$  is invertible.

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Signature is the number of positive eigenvalues minus the number of negative eigenvalues.

#### **Commuting case, again** The spectrum of

$$\mathcal{D}_{\lambda}(r_{1}, r_{2}, r_{3}) = \begin{bmatrix} r_{3} - \lambda_{3} & r_{1} - \lambda_{1} - i(r_{2} - \lambda_{2}) \\ r_{1} - \lambda_{1} + i(r_{2} - \lambda_{2}) & -r_{3} + \lambda_{3} \end{bmatrix}$$

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Thus

$$\lambda \in \Lambda(r_1, r_2, r_3)$$
 or  $\operatorname{Ind}_{\lambda}(\mathcal{D}_{\lambda}(r_1, r_2, r_3)) = 0$ .

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In the commuting case,

$$\operatorname{Ind}_{\lambda}(X, Y, Z) = 0$$

(or is undefined). From this one shows: for a given X, Y and Z there is some C so that

$$|\lambda| > C \implies \operatorname{Ind}_{\lambda}(X, Y, Z) = 0.$$

#### Theorem

If  $\operatorname{Ind}_{\lambda}(X, Y, Z) \neq \operatorname{Ind}_{\mu}(X, Y, Z)$  and  $\lambda_t$  is any curve between  $\lambda$  and  $\mu$  then there is at least one value of s so that

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#### Corollary

If for any  $\lambda$  we have  $\operatorname{Ind}_{\lambda}(X, Y, Z) \neq 0$  then  $\Lambda(X, Y, Z)$  separates  $\lambda$  from  $\infty$ .





$$(\omega = e^{\frac{2\pi i}{n}})$$

# **Fuzzy Tori** Consider $U = \begin{bmatrix} 0 & & & 1 \\ 1 & 0 & & \\ & \ddots & \ddots & \\ & & 1 & 0 \\ & & & & 1 & 0 \end{bmatrix}, \quad V = \begin{bmatrix} \omega & & & \\ & \omega^2 & & \\ & & \ddots & & \\ & & & \omega^{n-1} & \\ & & & & & 1 \end{bmatrix},$ (1) $(\omega = e^{\frac{2\pi i}{n}})$ which represent an embedding of a fuzzy torus into $\mathbb{R}^4$ .

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## **Fuzzy Tori**

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$$A_{n} = \frac{1}{2} \left( R + \frac{r}{2}U^{*} + \frac{r}{2}U \right) V^{*} + \frac{1}{2}V \left( R + \frac{r}{2}U^{*} + \frac{r}{2}U \right)$$
$$B_{n} = \frac{i}{2} \left( R + \frac{r}{2}U^{*} + \frac{r}{2}U \right) V^{*} - \frac{i}{2}V \left( R + \frac{r}{2}U^{*} + \frac{r}{2}U \right)$$
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Good choices for the radii are R = 0.8 and r = 0.4.

# **Emergent tori, showing** *K***-theory**



**Figure :** Left: approximation of part of a fuzzy torus, n = 4. Right: Same matrices, pseudospectrum cut open and showing some index values.



**Figure :** Now n = 6.



**Figure :** Now n = 10.



**Figure :** Now n = 20.



**Figure :** Now *n* = 50.

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For a (d-1)-dimensional physical system, the matrix model is then

 $(X_1, \ldots, X_{d-1}, H)$ 

where the  $X_j$  are commuting matrices for position, and H is the Hamiltonian, corresponding to energy.

# **Topological insulators**

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In the case d = 3, there is a simple model of what is called a Chern insulator. A homologically non-trivial surface emerged, as  $\Lambda(X, Y, H)$ .



**Figure :** Part of the pseudospectrum of a Chern insulator on an 18-by-18 lattice with no disorder. The vertical axis is energy, the others are position. The crosses at the Fermi level (energy zero) indicate nontrivial index, while the circles indicate trivial index. Reproduced from [3].



**Figure :** Part of the pseudospectrum of a Chern insulator on an 18-by-18 lattice with disorder sufficient to half close the gap in the bulk spectrum. Reproduced from [3].

# A finite volume index theorem

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#### Theorem (L. & Schulz-Baldes)

Suppose  $H = H^*$  is bounded on  $\mathbb{H} = \ell^2(\mathbb{Z}^2) \otimes \mathbb{C}^r$  is invertible and local. Specifically assume

 $\delta = \max\left( \left\| [\textbf{\textit{H}},\textbf{\textit{X}}] \right\|, \left\| [\textbf{\textit{H}},\textbf{\textit{Y}}] \right\| \right)$ 

is finite and  $\sigma(H) \cap (-\Delta, \Delta) = \emptyset$ . There exists  $L_0$  so that for  $L \ge L_0$  we have

$$\operatorname{Ind}_{\lambda}(X_{L}, Y_{L}, H_{L}) = \operatorname{Ind}_{P\mathbb{H}}\left(P\frac{X+iY}{|X+iY|}P\right)$$

where  $X_L$ ,  $Y_L$  and  $H_L$  are formed by restriction to  $\mathbb{H} = \ell^2([-L, L] \times [-L, L]) \otimes \mathbb{C}^r$ , with  $H_L$  given Dirichlet boundary conditions, and where P is the spectral projection for Hcorresponding to  $[0, \infty)$ .

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Also have results for *H* invertible and local on  $\mathbb{H} = \ell^2(\mathbb{Z}^{2d}) \otimes \mathbb{C}^r$ , etc.

Used the following to understand the boundary map.

$$0 \longrightarrow J \longrightarrow C_{\delta}(\mathbb{D}_d) \longrightarrow C_{\delta}(S^{d-1}) \longrightarrow 0$$

where

$$C_{\delta}(\mathbb{D}) = C_{1}^{*} \left\langle x_{1}, \dots, x_{d} \middle| \begin{array}{c} x_{j}^{*} = x_{j} \\ \|x_{j}x_{k} - x_{k}x_{j}\| \leq \delta \\ \sum x_{j}^{2} \leq 1 \end{array} \right\rangle$$

is a "soft *d*-ball" which is projective (L. & Shulman).

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about which we know less.

The usual formulas for mapping the *d*-ball onto the *d*-sphere lead to a coboundary map for

$$0 \rightarrow I \rightarrow A \rightarrow B \rightarrow 0$$

in term of "fuzzy" or "soft" spheres:

$$C_{\delta'}(S^{d-1}) o B \quad \hookrightarrow \quad C_{\delta}(S^d) o ilde{J}$$
Following the ideas in cohomotopy:



# References

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