# Emergent topology of insulators 

Noncommutative index theory Warsaw

Terry A. Loring

October, 2016

## Strings and D-branes

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Shown: fuzzy sphere, fuzzy torus.

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D_{(x, y, z)}=\left[\begin{array}{cc}
(X-x) & (Y-y)-i(Z-z) \\
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They call this the (shifted) Dirac operator.
It's square is almost the Laplace operator:

$$
D_{\mathbf{0}}^{2}=\left[\begin{array}{cc}
X^{2}+Y^{2}+Z^{2} & 0 \\
0 & X^{2}+Y^{2}+Z^{2}
\end{array}\right]+\left[\begin{array}{cc}
i[Y, Z] & {[X, Y]-i[X, Z]} \\
-[X, Y]-i[X, Z] & -i[Y, Z]
\end{array}\right]
$$

## Clifford spectrum - zero modes

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So the emergent surface is what mathematicians call the Joint Clifford Spectrum.

## Definition

For Hermitian operators $X, Y, Z$ the Clifford spectrum is the set

$$
\Lambda(X, Y, Z)=\left\{\lambda \in \mathbb{R}^{d} \mid D_{\lambda}(X, Y, Z) \text { is singular }\right\}
$$

## The unfuzzy case

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We have also the Laplacian:

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\Delta_{\lambda}(X, Y, Z)=\left(X-\lambda_{1}\right)^{2}+\left(Y-\lambda_{2}\right)^{2}+\left(Z-\lambda_{3}\right)^{2}
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In the case of commuting hermitian $X, Y$ and $Z$,
$\begin{aligned} & D_{\lambda}(X, Y, Z) \text { is singular } \\ \Longleftrightarrow & \Delta_{\lambda}(X, Y, Z) \text { is singular } \\ \Longleftrightarrow & \exists \mathbf{v} \text { a unit vector with } X \mathbf{v}=\lambda_{1} \mathbf{v}, Y \mathbf{v}=\lambda_{2} \mathbf{v}, Z \mathbf{v}=\lambda_{3} \mathbf{v}\end{aligned}$

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This means $\Lambda(X, Y, Z)$ is a finite set.

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& X_{n}=\frac{1}{n}\left[\begin{array}{lllll}
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& n-1 & & & \\
& & \ddots & & \\
& & & 1-n & \\
& & & & -n
\end{array}\right], \\
& Y_{n}=\frac{1}{2 n}\left(T_{n}^{*}+T_{n}\right), \quad Z_{n}=\frac{i}{2 n}\left(T_{n}^{*}-T_{n}\right)
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$$

where we set $N_{n}=n(n+1)$ and
$T_{n}=\left[\begin{array}{ccccc}0 & \sqrt{N_{n}-(n-1) n} & & & \\ & 0 & \sqrt{N_{n}-(n-2)(n-1)} & & \\ & & \ddots & \ddots & \\ & & & 0 & \sqrt{N_{n}-(-n)(1-n)}\end{array}\right]$

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The second example is $X_{n}, \tilde{Y}_{n}, \tilde{Z}_{n}$ where

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$$ and

$$
S_{n}=\left[\begin{array}{ccccc}
0 & 1 & & & \\
& 0 & 1 & & \\
& & \ddots & \ddots & \\
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For Hermitian operators $X, Y, Z$ and $\epsilon \geq 0$, define the Clifford $\epsilon$-pseudospectrum to be the set

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with the convention that $R$ singular means $\left\|R^{-1}\right\|^{-1}=0$.

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Notice

$$
\left\|\not D_{\lambda}(X, Y, Z)^{-1}\right\|^{-1}=\left|\operatorname{eig}_{\min }\left(D_{\lambda}(X, Y, Z)\right)\right|
$$

called the gap at $\lambda$.


Figure: Left: approximation of part of the Clifford spectrum for the usual fuzzy sphere, spin=2. Right: Same spin but with modified $Y$ and $Z$.


Figure : Now spin=20.


Figure: Now spin=200.


Figure: Now spin=2000.

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The local index for the finite system $(X, Y, Z)$ is then

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\operatorname{Ind}_{\lambda}(X, Y, Z)=\frac{1}{2} \operatorname{Sig}\left(D_{\lambda}(X, Y, Z)\right)
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defined at every point where $D_{\lambda}$ is invertible.

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Signature is the number of positive eigenvalues minus the number of negative eigenvalues.

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The spectrum of

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Thus

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\lambda \in \Lambda\left(r_{1}, r_{2}, r_{3}\right) \quad \text { or } \quad \operatorname{Ind}_{\lambda}\left(D_{\lambda}\left(r_{1}, r_{2}, r_{3}\right)\right)=0
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(or is undefined). From this one shows: for a given $X, Y$ and $Z$ there is some $C$ so that

$$
|\lambda|>C \Longrightarrow \operatorname{Ind}_{\lambda}(X, Y, Z)=0
$$

## Theorem

If $\operatorname{Ind}_{\lambda}(X, Y, Z) \neq \operatorname{Ind}_{\mu}(X, Y, Z)$ and $\lambda_{t}$ is any curve between $\lambda$ and $\mu$ then there is at least one value of $s$ so that

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\lambda_{s} \notin \Lambda(X, Y, Z) .
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## Corollary

If for any $\lambda$ we have $\operatorname{Ind}_{\lambda}(X, Y, Z) \neq 0$ then $\Lambda(X, Y, Z)$ separates $\lambda$ from $\infty$.

## Fuzzy Tori

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Consider

$$
U=\left[\begin{array}{ccccc}
0 & & & & 1 \\
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& \ddots & \ddots & & \\
& & 1 & 0 & \\
& & & 1 & 0
\end{array}\right], \quad V=\left[\begin{array}{lllll}
\omega & & & & \\
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& & \ddots & & \\
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( $\omega=e^{\frac{2 \pi i}{n}}$ ) which represent an embedding of a fuzzy torus into $\mathbb{R}^{4}$. We seek an embedding into $\mathbb{R}^{3}$. Set

$$
\begin{gathered}
A_{n}=\frac{1}{2}\left(R+\frac{r}{2} U^{*}+\frac{r}{2} U\right) V^{*}+\frac{1}{2} V\left(R+\frac{r}{2} U^{*}+\frac{r}{2} U\right) \\
B_{n}=\frac{i}{2}\left(R+\frac{r}{2} U^{*}+\frac{r}{2} U\right) V^{*}-\frac{i}{2} V\left(R+\frac{r}{2} U^{*}+\frac{r}{2} U\right) \\
C_{n}=\frac{r i}{2} U^{*}-U \frac{r i}{2} .
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Good choices for the radii are $R=0.8$ and $r=0.4$.

## Emergent tori, showing K-theory




Figure : Left: approximation of part of a fuzzy torus, $n=4$. Right: Same matrices, pseudospectrum cut open and showing some index values.


Figure : Now $n=6$.


Figure: Now $n=10$.


Figure: Now $n=20$.


Figure: Now $n=50$.

## Insulators

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In finite models of condensed matter, many use periodic boundary conditions.
Edge modes are important; I prefer open boundaries.
For a ( $d-1$ )-dimensional physical system, the matrix model is then

$$
\left(X_{1}, \ldots, X_{d-1}, H\right)
$$

where the $X_{j}$ are commuting matrices for position, and $H$ is the Hamiltonian, corresponding to energy.

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In the case $d=3$, there is a simple model of what is called a Chern insulator. A homologically non-trivial surface emerged, as $\Lambda(X, Y, H)$.


Figure : Part of the pseudospectrum of a Chern insulator on an 18-by-18 lattice with no disorder. The vertical axis is energy, the others are position. The crosses at the Fermi level (energy zero) indicate nontrivial index, while the circles indicate trivial index. Reproduced from [3].


Figure: Part of the pseudospectrum of a Chern insulator on an 18-by-18 lattice with disorder sufficient to half close the gap in the bulk spectrum. Reproduced from [3].

## A finite volume index theorem

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## Theorem (L. \& Schulz-Baldes)

Suppose $H=H^{*}$ is bounded on $\mathbb{H}=\ell^{2}\left(\mathbb{Z}^{2}\right) \otimes C^{r}$ is invertible and local. Specifically assume

$$
\delta=\max (\|[H, X]\|,\|[H, Y]\|)
$$

is finite and $\sigma(H) \cap(-\Delta, \Delta)=\varnothing$. There exists $L_{0}$ so that for $L \geq L_{0}$ we have

$$
\operatorname{Ind}_{\lambda}\left(X_{L}, Y_{L}, H_{L}\right)=\operatorname{Ind}_{P H}\left(P \frac{X+i Y}{|X+i Y|} P\right)
$$

where $X_{L}, Y_{L}$ and $H_{L}$ are formed by restriction to $\mathbb{H}=\ell^{2}([-L, L] \times[-L, L]) \otimes \mathbb{C}^{r}$, with $H_{L}$ given Dirichlet boundary conditions, and where $P$ is the spectral projection for $H$ corresponding to $[0, \infty)$.

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Also have results for $H$ invertible and local on $\mathbb{H}=\ell^{2}\left(\mathbb{Z}^{2 d}\right) \otimes \mathbb{C}^{r}$, etc.

Used the following to understand the boundary map.

$$
0 \longrightarrow J \longrightarrow C_{\delta}\left(\mathbb{D}_{d}\right) \longrightarrow C_{\delta}\left(S^{d-1}\right) \longrightarrow 0
$$

where

$$
C_{\delta}(\mathbb{D})=C_{1}^{*}\left\langle\begin{array}{l|c}
x_{1}, \ldots, x_{d} & x_{j}^{*}=x_{j} \\
\left\|x_{j} x_{k}-x_{k} x_{j}\right\| \leq \delta \\
\sum x_{j}^{2} \leq 1
\end{array}\right\rangle
$$

is a "soft d-ball" which is projective (L. \& Shulman).

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x_{j}^{*}=x_{j} \\
\left\|x_{j} x_{k}-x_{k} x_{j}\right\| \leq \delta \\
1-\delta \leq \sum x_{j}^{2} \leq 1
\end{array}
\end{array}\right\rangle
$$

about which we know less.

The usual formulas for mapping the $d$-ball onto the $d$-sphere lead to a coboundary map for

$$
0 \rightarrow I \rightarrow A \rightarrow B \rightarrow 0
$$

in term of "fuzzy" or "soft" spheres:

$$
C_{\delta^{\prime}}\left(S^{d-1}\right) \rightarrow B \quad \uparrow \quad C_{\delta}\left(S^{d}\right) \rightarrow \tilde{J}
$$

Following the ideas in cohomotopy:


## References

T
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䡒 Joanna L. Karczmarek and Ken Huai-Che Yeh. Noncommutative spaces and matrix embeddings on flat $\mathbb{R}^{2 n+1}$. Journal of High Energy Physics, 2015(11):1-15, 2015.
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