

Asymptotics of solutions in nonlocal spatially continuous parabolic Anderson-type models with Poissonian interaction: annealed and quenched behaviour

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Introduction

We will be interested in the long-time behaviour of solutions to **parabolic Anderson problem** on \mathbb{R}^d (abbr. PAM) driven by the operator L :

$$\begin{aligned}u_t^\omega(t, x) &= L_x u^\omega(t, x) - V^\omega(x) u^\omega(t, x), \quad t > 0, \\u^\omega(0, x) &= 1,\end{aligned}$$

where

- L is the generator of a Lévy process on \mathbb{R}^d , (L_x indicates that the operator acts on the space variable x),
- V^ω is a random potential of **Poissonian type** over some probability space $(\Omega, \mathcal{M}, \mathbb{P})$, corresponding to random impurities in the medium.

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One is interested in the asymptotics of

- $\mathbb{E}[u^\omega(t, x)]$ (the annealed behaviour of u^ω),
- $u^\omega(t, x)$, \mathbb{P} -a.s. (the quenched behaviour of u^ω),

as $t \rightarrow \infty$.

Settings – probabilistic

- $X = (X_t)_{t \geq 0}$ – **symmetric Lévy process** in \mathbb{R}^d , $d \geq 1$
- $\mathbf{P}_x, \mathbf{E}_x$ – measure and expected value of X_t starting at $x \in \mathbb{R}^d$
- characteristic function: $\mathbf{E}^0 e^{i\xi \cdot X_t} = e^{-t\psi(\xi)}$, $t > 0$, where

$$\psi(\xi) = \xi \cdot A\xi + \int_{\mathbb{R}^d \setminus \{0\}} (1 - \cos(\xi \cdot z)) \nu(dz)$$

A – symmetric non-negative definite $d \times d$ matrix (*Gaussian coefficient*)
 $\nu(dz)$ – symmetric m. on $\mathbb{R}^d \setminus \{0\}$, $\int (1 \wedge |z|^2) \nu(dz) < \infty$ (*Lévy measure*)
(we need some also some mild assumptions concerning ψ).

ψ is called the **Lévy-Khinchine exponent** of the process X

- generator: for $\varphi \in C_c^2(\mathbb{R}^d)$,

$$L\varphi(x) = \lim_{t \rightarrow 0} \frac{\mathbf{E}_x[\varphi(X_t)] - \varphi(x)}{t},$$

the limit taken in $L^2(\mathbb{R}^d)$.

Settings – analytic

The generator of the process can be expressed analytically:

$$\widehat{L}f(\xi) = -\psi(\xi)\widehat{f}(\xi), \quad f \in D(L) := \left\{ f \in L^2(\mathbb{R}^d) : \psi\widehat{f} \in L^2(\mathbb{R}^d) \right\};$$

for $\varphi \in C_c^2(\mathbb{R}^d)$

$$L\varphi(x) = \sum_{j,k=1}^d a_{jk} \frac{\partial^2 \varphi}{\partial x_j \partial x_k}(x) + \int_{\mathbb{R}^d} (\varphi(x+y) - \varphi(x) - \nabla \varphi(x) \cdot y \mathbf{1}_{\{|y| \leq 1\}}) \nu(dy).$$

Example: symmetric stable processes

$X^{(\alpha)} = (X_t^{(\alpha)})_{t \geq 0}$ – symmetric α -stable process with L-K exponent $\psi^{(\alpha)}$:

$$\psi^{(\alpha)}(\xi) = \begin{cases} \int_0^\infty \int_{S^{d-1}} \frac{1 - \cos(\xi \cdot rz)}{r^{1+\alpha}} n(dz) dr & \text{when } \alpha \in (0, 2) \\ \xi \cdot A \xi & \text{when } \alpha = 2 \end{cases}$$

n – symmetric finite measure on unit sphere S^{d-1}

$A = (a_{ij})_{1 \leq i, j \leq d}$ – symmetric nonnegative definite matrix

When n is the uniform distribution on S^{d-1} for $\alpha \in (0, 2)$ or $A \equiv a \text{Id}$ with some $a > 0$ for $\alpha = 2$, then the process is called *isotropic α -stable*.

We assume the nondegeneracy condition $\inf_{|\xi|=1} \psi^{(\alpha)}(\xi) > 0$.

For the isotropic α -stable process, $\alpha \in (0, 2)$, the generator is the fractional Laplacian: $L = -(-\Delta)^{\alpha/2}$.

Other examples

- Mixture of stable processes: $\psi(\xi) = a_0|\xi|^2 + \sum_{i=1}^n a_i|\xi|^{\alpha_i}$,
 $a_0 \geq 0$, $a_i > 0$ for $i = 1, 2, \dots, n$, $\alpha_i \in (0, 2)$.
- Geometric stable process with Gaussian component:
 $\psi(\xi) = \xi \cdot A\xi + \log(1 + |\xi|^\delta)$ with nondegenerate A and
 $\delta \in (0, 2)$.
- Relativistic α -stable processes: $\psi(\xi) = (|\xi|^2 + m^{2/\alpha})^{\alpha/2} - m$,
 $\alpha \in (0, 2)$, $m > 0$.
- Layered stable processes, tempered stable processes, Lamperti
stable processes, truncated stable processes...

Principal eigenvalues

Notation:

$\lambda_1^{(\alpha)}(U)$ – principal eigenvalue of the process $(X_t^{(\alpha)})_{t \geq 0}$ killed on leaving an open and bounded set $U \subset \mathbb{R}^d$

$(0 < \lambda_1^{(\alpha)}(U) := \inf \text{spec}(-L_U^{(\alpha)}))$

$\lambda_1^{BM}(U)$ – the case of Brownian motion ($A \equiv Id$)

Poisson point process and Poissonian random field

Poisson random measure with intensity ρdx , $\rho > 0$, over $(\Omega, \mathcal{M}, \mathbb{P})$ (the Poisson cloud)

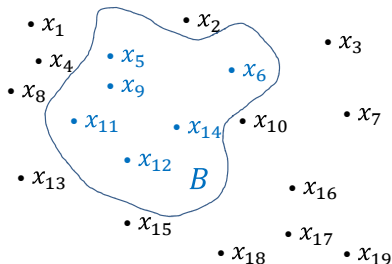
$$\mu^\omega(dy) = \sum_{i=1}^{\infty} \delta_{x_i}(dy)$$

$$x_i \sim \omega \in \Omega$$

(δ_x – the Dirac point measure at x)

$(\Omega, \mathcal{M}, \mathbb{P})$ – new probability space,
different from that of the process X
(full **stochastic independence**)

\mathbb{E} – expected value w.r.t. \mathbb{P}



For each Borel set $B \in \mathcal{B}(\mathbb{R}^d)$ with $0 < |B| < \infty$ and $n = 0, 1, 2, \dots$

$$\mathbb{P}(\mu^\omega(B) = n) = \exp(-\rho|B|) \frac{(\rho|B|)^n}{n!} \quad [\text{Poisson distribution with parameter } \rho|B|],$$

and for disjoint sets A, B the clouds over A, B are independent.

Poisson point process and Poissonian random field

Poisson random potential

- $W : \mathbb{R}^d \rightarrow [0, \infty)$, $W \in L^\infty(\mathbb{R}^d)$, with bounded support (*potential profile*)
- define

$$V^\omega(x) = \int_{\mathbb{R}^d} W(x-y)\mu^\omega(dy) = \sum_{x_i \in \omega} W(x-x_i), \quad x \in \mathbb{R}^d, \quad \omega \in \Omega$$

(*Poissonian potential*)

Poisson point process and Poissonian random field

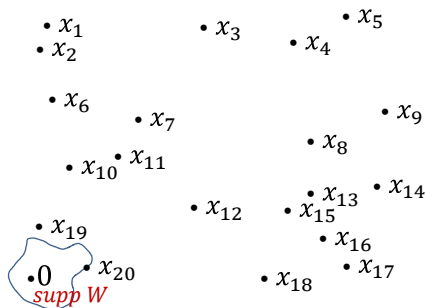
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(*Poissonian potential*)

e.g. $W(x) := \mathbf{1}_E(x)$
for some $E \subset \mathcal{B}(\mathbb{R}^d)$



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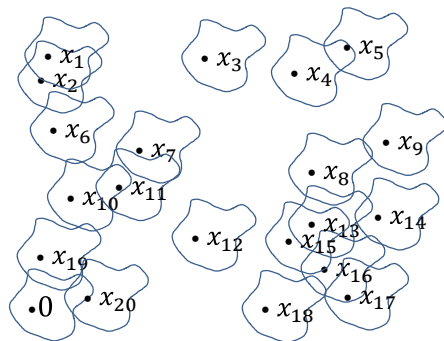
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$$\begin{aligned} V^\omega(x) &= \sum_{x_i \in \omega} \mathbf{1}_E(x-x_i) \\ &= \sum_{x_i \in \omega} \mathbf{1}_{E+x_i}(x) \end{aligned}$$



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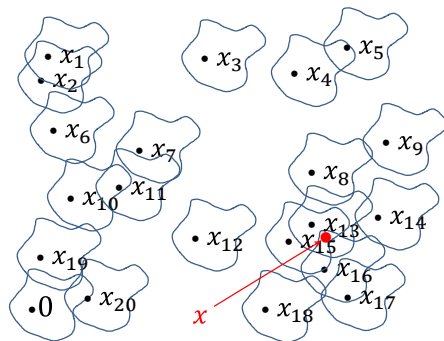
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$$V^\omega(x) = \sum_{i \in \{13,15,16\}} \mathbf{1}_{E+x_i}(x) = 3$$



Lévy process in Poissonian environment

For $\omega \in \Omega$, $B \in \mathcal{B}(\mathbb{R}^d)$, $x \in \mathbb{R}^d$ and $t > 0$ consider

$$\mathbf{P}_x^\omega(X_t \in B) = \mathbf{E}_x[e^{-\int_0^t v^\omega(X_s) ds}; X_t \in B] = \mathbf{E}_x[e^{-\sum_{x_i(\omega)} \int_0^t W(X_s - x_i) ds}; X_t \in B]$$

Survival probability of the Lévy process in Poissonian environment

$$u^\omega(t, x) := \mathbf{E}_x[e^{-\int_0^t v^\omega(X_s) ds}], \quad t > 0, \quad x \in \mathbb{R}^d$$

Lévy process in Poissonian environment

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Survival probability of the Lévy process in Poissonian environment

$$u^\omega(t, x) := \mathbf{E}_x[e^{-\int_0^t V^\omega(X_s) ds}], \quad t > 0, \quad x \in \mathbb{R}^d$$

The function u^ω is the solution of (Feynman-Kac formula)

$$\frac{\partial u}{\partial t} = -H^\omega u, \quad u(0, x) \equiv 1,$$

where

$$H^\omega = -L + V^\omega$$

Brownian motion in Poissonian medium: $H^\omega = -\Delta + V^\omega$

Annealed behaviour [Donsker-Varadhan, CPAM, 1975]

$$u^\omega(t, x) = \mathbb{E} \mathbf{E}_x \left[e^{-\int_0^t V^\omega(X_s) ds} \right] = \exp \left(-C t^{d/(d+2)} (1 + o(1)) \right), \quad \text{as } t \rightarrow \infty,$$

with

$$C = (\rho \omega_d)^{\frac{2}{d+2}} \left(\frac{d+2}{2} \right) \left(\frac{2\lambda_1^{BM}(B(0,1))}{d} \right)^{\frac{d}{d+2}}$$

ω_d – the volume of a unit ball

Quenched behaviour [Sznitman, PTRF, 1993]

$$u^\omega(t, x) = \exp \left(-C \frac{t}{(\log t)^{2/d}} (1 + o(1)) \right), \quad \text{as } t \rightarrow \infty, \quad \mathbb{P} - \text{a.s.},$$

with

$$C = \left(\frac{\rho \omega_d}{d} \right)^{\frac{2}{d}} \lambda_1^{BM}(B(0,1))$$

Quenched behaviour obtained from annealed one [R. Fukushima, ECP, 2009]

Lévy process with jumps: $H^\omega = -L + V^\omega$

Annealed behaviour [Okura, Z. Wahrscheinlichkeitstheorie verw. Gebiete, 1981]

If

$$(A) \quad \psi(\xi) = \psi^{(\alpha)}(\xi) + o(|\xi|^\alpha), \quad |\xi| \rightarrow 0,$$

for some stable process $X^{(\alpha)}$ with $\alpha \in (0, 2]$, and

$$\frac{\psi(\xi)}{(\log |\xi|)^2} \rightarrow \infty \quad \text{as } |\xi| \rightarrow \infty,$$

then

$$u^\omega(t, x) = \exp\left(-Ct^{d/(d+\alpha)}(1 + o(1))\right), \quad \text{as } t \rightarrow \infty,$$

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$$\lambda_{(\alpha)} = \inf_{U \text{ open}, |U|=\omega_d} \lambda_1^{(\alpha)}(U).$$

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What about the quenched behaviour for jump processes?

The quenched behaviour. The upper bound.

(U) There exist constants $C > 0$, $r_0 \geq 1$, $t_1 \geq t_0 \vee 1$ and a nonincreasing function $F : [0, \infty) \rightarrow [0, \infty)$ such that $\lim_{r \rightarrow \infty} F(r) = 0$, for which

$$P_0(|X_t| \geq r) \leq C t (F(r) \vee e^{-r}), \quad r \geq r_0 \vee 2t, \quad t \geq t_1.$$

Theorem (KK and KPP, 2016)

Suppose that X is a Lévy process satisfying **(A)** and **(U)**. Then there exist a function $h : \mathbb{R}_+ \rightarrow \mathbb{R}_+$ such that for every $x \in \mathbb{R}^d$

$$\limsup_{t \rightarrow \infty} \frac{\log u^\omega(t, x) - \frac{d}{2} \log h(t)}{g(t)} \leq - \left(\frac{\rho}{d}\right)^{\alpha/d} \lambda_1^{(\alpha)}, \quad \mathbb{Q} - a.s.,$$

where $g(t) = t/(\log h(t))^{\alpha/d}$.

Recall that $\lambda_1^{(\alpha)} = \inf_{U \text{ open}, |U| = \omega_d} \lambda_1^{(\alpha)}(U)$.

The rate function

What are the **correction term** h and the **rate function** g ?

For fixed $\kappa > 0$, $\alpha \in (0, 2]$ let $f_{\alpha, \kappa} : [1, \infty) \rightarrow [0, \infty)$ be given by

$$f_{\alpha, \kappa}(r) = \left((r \wedge |\log(1 \wedge F(r))|) + \frac{d}{2} \log r \right) \left(\frac{d \log r}{\kappa} \right)^{\frac{\alpha}{d}}.$$

Let $h_{\alpha, \kappa} = f_{\alpha, \kappa}^{-1}$.

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It is well defined since $f_{\alpha, \kappa}(1) = 0$, $\lim_{r \rightarrow \infty} f_{\alpha, \kappa}(r) = \infty$, moreover $f_{\alpha, \kappa}(r)$ is **strictly increasing** in r .

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$$t \left(\frac{\kappa}{d \log h_{\alpha, \kappa}(t)} \right)^{\frac{\alpha}{d}} = \left(h_{\alpha, \kappa}(t) \wedge |\log(1 \wedge F(h_{\alpha, \kappa}(t)))| \right) + \frac{d}{2} \log h_{\alpha, \kappa}(t).$$

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We take $h = h_{\alpha, \kappa_0}$ with $\kappa_0 = \rho(\lambda_1^{(\alpha)})^{d/\alpha}$.

The upper bound – special cases

(U) There exist constants $C > 0$, $r_0 \geq 1$, $t_1 \geq t_0 \vee 1$ and a nonincreasing function $F : [0, \infty) \rightarrow [0, \infty)$ such that $\lim_{r \rightarrow \infty} F(r) = 0$, for which

$$P_0(|X_t| \geq r) \leq C t (F(r) \vee e^{-r}), \quad r \geq r_0 \vee 2t, \quad t \geq t_1.$$

Corollary

Suppose that the function F appearing in **(U)** satisfies $\lim_{t \rightarrow \infty} \frac{|\log F(r)|}{\log r} = +\infty$.

Then $\lim_{t \rightarrow \infty} \frac{h(t)}{g(t)} = 0$, so that for every $x \in \mathbb{R}^d$

$$\limsup_{t \rightarrow \infty} \frac{\log u^\omega(t, x)}{g(t)} \leq - \left(\frac{\rho}{d}\right)^{\alpha/d} \lambda_1^{(\alpha)}, \quad \mathbb{P} - a.s..$$

For **stable processes with $\alpha \in (0, 2)$** and for mixtures of stable processes with $\alpha_i \in (0, 2)$, one has $g(t) = ct^{\frac{d}{d+\alpha}}$. For **the relativistic process and for the Brownian motion**, $g(t) = ct/(\log t)^{2/d}$.

Quenched behaviour. The lower bound

Let $B_R =: [-R, R]^d$ and let $p^{B_R}(t, x, y)$ be the Dirichlet kernel of the process X in B_R . Let

$$G(R) := \inf_{|y| \leq R/2} p^{B_R}(1, 0, y).$$

Theorem (KK and KPP, 2016)

Let X be a Lévy process with exponent ψ satisfying **(A)** and let the potential profile W be of finite range. Then for any $x \in \mathbb{R}^d$, \mathbb{Q} -almost surely

$$\begin{aligned} \liminf_{t \rightarrow \infty} \frac{\log u^\omega(t, x) + 3d \log h(t) + \left| \log G \left(\frac{2h(t)}{(\log h(t))^{\frac{\alpha}{d} + 2}} \right) \right|}{g(t)} \\ \geq - \left(\frac{\omega_d \rho}{d} \right)^{\alpha/d} \lambda_1^{(\alpha)}(B(0, 1)), \end{aligned}$$

where $h(t)$, $g(t)$ are as before, and $\lambda_1^{(\alpha)}(B(0, 1))$ is the principal Dirichlet eigenvalue of $B(0, 1)$ for the α -stable process from **(A)**.

Quenched behaviour for isotropic stable process

Theorem (KK and KPP, 2016)

If

$$\psi(\xi) = c|\xi|^\alpha, \quad \xi \in \mathbb{R}^d,$$

for some $\alpha \in (0, 2)$ then for any $x \in \mathbb{R}^d$, \mathbb{P} -almost surely,

$$\exp\left(-C_L t^{d/(d+\alpha)}(1+o(1))\right) \leq u^\omega(t, x) \leq \exp\left(-C_U t^{d/(d+\alpha)}(1+o(1))\right)$$

as $t \rightarrow \infty$, with

$$C_L = \frac{4\alpha + 9d}{2} \left(\frac{2}{d+2\alpha}\right)^{\frac{d}{d+\alpha}} \left(\frac{\rho\omega_d}{d}\right)^{\frac{\alpha}{d+\alpha}} \left(\lambda_1^{(\alpha)}(B(0,1))\right)^{\frac{d}{d+\alpha}}$$

and

$$C_U = \alpha \left(\frac{2}{d+2\alpha}\right)^{\frac{d}{d+\alpha}} \left(\frac{\rho\omega_d}{d}\right)^{\frac{\alpha}{d+\alpha}} \left(\lambda_1^{(\alpha)}(B(0,1))\right)^{\frac{d}{d+\alpha}}.$$

Quenched behaviour for relativistic stable processes

Theorem (KK and KPP, 2016)

Let X_t be the α -stable relativistic process, i.e. the process with Lévy-Khinchine exponent $\psi(\xi) = (|\xi|^2 + m^{2/\alpha})^{\alpha/2} - m$, $m > 0$, $\alpha \in (0, 2)$. Then for any $x \in \mathbb{R}^d$, \mathbb{P} -almost surely

$$u^\omega(t, x) = \exp(-C_{m,\alpha}(t/(\log t)^{2/d})(1 + o(1))),$$

where the constant $C_{m,\alpha}$ is given by

$$C_{m,\alpha} = \frac{\alpha}{2} m^{1-\frac{2}{\alpha}} \alpha_1^{(2)}(B(0, 1)) \left(\frac{\rho\omega_d}{d}\right)^{\frac{d}{2}}.$$









- Observe that when $\alpha \rightarrow 2$, then the constant tends to $\alpha_1^{BM}(B(0, 1)) \left(\frac{\rho\omega_d}{d}\right)^{d/2}$, which is the constant for the Brownian motion.

Annealed and quenched asymptotics

Quenched asymptotics depends on the rate of decay of the Lévy measure at ∞ .

$\nu(x), x \rightarrow \infty$	parameters	quenched rate	annealed rate
$\frac{C}{ x ^{d+\alpha}}$	$\alpha \in (0, 2)$	$t^{\frac{d}{d+\alpha}}$	$t^{\frac{d}{d+\alpha}}$
$\frac{C}{ x ^{d+\delta}}$	$\delta > 2$	$t^{\frac{d}{d+2}}$	$t^{\frac{d}{d+2}}$
$\frac{C}{e^{\theta(\log x)^\beta}}$	$\theta > 0$ $\beta > 1$	$t^{\frac{d\beta}{d\beta+2}}$	$t^{\frac{d}{d+2}}$
$\frac{C}{e^{\theta x ^\beta}}$	$\theta > 0$ $\beta > 0$	$\frac{t}{(\log t)^{2/d}}$	$t^{\frac{d}{d+2}}$
0		$\frac{t}{(\log t)^{2/d}}$	$t^{\frac{d}{d+2}}$
BM		$\frac{t}{(\log t)^{2/d}}$	$t^{\frac{d}{d+2}}$

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