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in collaboration with
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# Phase structure of Causal Dynamical Triangulations model in 4D 

Singularities of general relativity and their quantum fate

29th June 2016

UNIWERSYTET
JAGIELLOŃSKI
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$\triangleleft$ But it could be renormalizable in a nonperturbative regime
$\diamond$ assymptotic safety idea (S. Weinberg)
$\diamond$ renormalization group flow can lead to a nonGaussian UV fixed point (M. Reuter et al.)
\& Lattice formulation would allow to study non-perturbative gravity
$\diamond$ we need a dynamical lattice (DT)
$\diamond$ UV fixed point should be associated with a 2nd order phase transition
$\diamond$ one should be able to reproduce semi-classical gravity (IR limit)
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## Outline

$\diamond C D T$
$\diamond$ Phase structure
$\diamond$ Phase transitions
$\diamond$ Transfer matrix method
$\triangleleft$ Phase structure revisited
$\diamond$ Phase transitions revisited
$\diamond$ Signature change
$\diamond$ Conclusions

## CDT

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¿ Classical mechanics: single trajectory of a particle resulting from $E$-L equations (Hamilton's principle)
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> Einstein's General Relativity: gravity defined through spacetime geometry
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$\diamond$ We will consider pure gravity model $(G)$ with positive cosmological $\quad S_{H E}=\frac{1}{16 \pi G} \int d^{4} x \sqrt{-g}(R-2 \Lambda)$ constant ( $\wedge$ )

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$\triangleleft$ CDT is formulated in a coordinate free way
$\diamond$ Three coupling constants: $k_{0}, k_{4}, \Delta$
४ After Wick's rotation: „random" geometry system

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Z=\int_{\text {trajectories }} D\left[g_{\mu \nu}\right] \exp \left(i S_{H E}\left[g_{\mu \nu}\right]\right)
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$\triangleleft$ The observable: 3-volume of spatial layers (foliation leaves of the global proper time): $V_{3}(t) \propto n_{t} \equiv N_{(4,1)}(t)$
$\triangleleft$ Initially three phases $(A, B, C)$ of various geometry were discovered
$\triangleleft$ Phase C (de Sitter phase) has good semi-classical properties (IR limit)
$\diamond$ Hausdorff dim.: 4 , spectral dim.: $2 \Rightarrow 4$
$\diamond$ Background geommetry is consistent with a 4-dim sphere $\Rightarrow$ Euclidean de Sitter universe (positive cosmol. const.)
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\begin{aligned}
\bar{n}_{t} \equiv\left\langle n_{t}\right\rangle & =\frac{3}{4} \tilde{V}_{4} \frac{1}{\tilde{A} \tilde{V}_{4}^{1 / 4}} \cos ^{3}\left(\frac{t-t_{0}}{\tilde{A} \tilde{V}_{4}^{1 / 4}}\right) \\
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\begin{gathered}
d s^{2}=d t^{2}+a^{2}(t) d \Omega_{3}^{2} \Rightarrow V_{3}(t) \propto a^{3}(t) \\
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S=-\frac{1}{24 \pi G} \int d t\left(\frac{\dot{V}_{3}(t)^{2}}{V_{3}(t)}+\mu V_{3}(t)^{1 / 3}-\lambda V_{3}(t)\right)
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$\triangleleft$ Phase C (de Sitter phase) continued
$\diamond$ The effective action for the $n_{t}$ observable ...
४ ... can be analyzed by looking at quantum fluctuations around the semiclassical solution
$\triangleleft$ The (inverse of) covariance matrix $P=C^{-1}$ provides information about second derivatives of the effective action
$\diamond$ The measured covariance matrix is consistent with MS action (with reversed overall sign)!

$$
\begin{aligned}
& \text { ed } \\
& \begin{array}{l}
Z=\sum_{T} \exp \left(-S_{R}[T]\right)=\sum_{\left\{n_{t}\right\} T\left[\left\{n_{n}\right\}\right]}^{\prime} \sum_{2}^{\prime} \exp \left(-S_{R}[T]\right) \\
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n_{t}=\left\langle n_{t}\right\rangle+\delta n_{t} \quad C_{t t^{\prime}} \equiv\left\langle\delta n_{t} \delta n_{t^{\prime}}\right\rangle
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S_{e f}=\frac{1}{\Gamma} \sum_{t}\left(\frac{\left(n_{t+1}-n_{t}\right)^{2}}{\left(n_{t}+n_{t+1}\right)}+\tilde{\mu} n_{t}^{1 / 3}-\tilde{\lambda} n_{t}\right) \\
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## Phase transitions

$\diamond$ To analyze phase transitions one needs to define a suitable order parameter OP (e.g. $N_{4}, N_{0}, \ldots$ )
$\triangleleft$ Here we look at OP conjugate to the varied coupling constant
২ (Pseudo)critical point is signaled by max. of susceptibility
$\diamond$ Two-states jumping of OP (double
 peak structure of measured histograms) may signal a 1st order transitions
$\triangleleft$ But one must be careful and check $N_{4} \rightarrow \infty$ limit
২ There exists a 2-nd order transition
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## Transfer matrix method

$\diamond$ The transfer matrix method enables one to measure the effective action directly
$\triangleleft$ CDT has by definition a transfer matrix parametrized by 3-dimensional spatial triangulations $T_{3}$
४ Local form of the effective action in Phase C suggests that a description by effective transfer matrix parametrized by spatial volume $n_{t}$ is also viable
« Measurement of the transfer matrix $=$ direct measurement of the effective Lagrangian

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४ Measurement of the transfer matrix $=$

$$
\begin{gathered}
S_{e f}=\frac{1}{\Gamma} \sum \sum_{t}\left(\frac{\left(n_{t+1}-n_{t}\right)^{2}}{\left(n_{t}+n_{t+1}\right)}+\tilde{\mu} n_{t}^{1 / 3}-\tilde{\lambda}_{n t}\right) \\
S_{e f}=\sum_{t} L_{e f}\left[n_{t}, n_{t+1}\right] \\
Z_{e f}=\sum_{\left\{n_{t}\right\}}\left\langle n_{t}\right| M_{e f}^{T}\left|n_{t}\right\rangle=t r M_{e f}^{T}
\end{gathered}
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$\diamond$ Gaussian kinetic term \& MS potential
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$\diamond$ Infinite Hausdorff dimension?
$\triangleleft$ Spectral dimension $>4$ and growing (to infinity ?) with growing volume
४ This suggests high connectivity between the building blocks
४ 4-volume is concentrated in short geodesic distance
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$\diamond$ whose boundary has topology of a 2-sphere (Euler characteristic = 2)
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 can expand in powers of $2 c_{0}(n-m) / \Gamma \ll 1$

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\langle n| M_{C}|m\rangle=N[n+m] \exp \left(-\frac{(m-n)^{2}}{\Gamma(n+m)}\right)
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## Conclusions

$\diamond$ Using the transfer matrix method a new D (bifurcation) phase was discovered JHEP1406(2014)034
$\diamond$ The geometry inside the new phase is very non-trivial and much different than inside the C (de Sitter) phase JHEP1508(2015)033
$\diamond$ The transition seems to be 2nd order JHEP1602(2016)144
$\diamond$ New phase transition may be related to signature change JHEP1508(2015)033
Prospects
$\diamond$ Different topology of spatial slices $\left(S^{3} \Rightarrow T^{3}\right)$
$\diamond$ Inclusion of matter (massless scalar fields)


## Thank You!



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