

Some relationships between potential theories of classical and hyperbolic Brownian motion

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- Hyperbolic distance formula

$$\cosh(d_{\mathbb{H}^n}(x, y)) = 1 + \frac{|x - y|^2}{2x_n y_n}.$$

- Volume element

$$dV_n = \frac{1}{x_n^n} dx_1 \dots dx_n$$

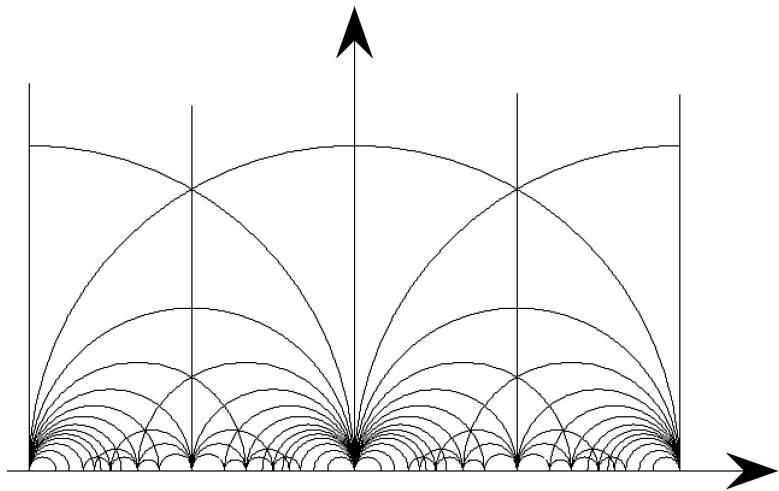
- Volume element

$$dV_n = \frac{1}{x_n^n} dx_1 \dots dx_n$$

- Laplace - Beltrami operator

$$\Delta = \operatorname{div}(\operatorname{grad}) = x_n^2 \sum_{k=1}^n \frac{\partial^2}{\partial x_k^2} - (n-2)x_n \frac{\partial}{\partial x_n}$$

Hyperbolic space \mathbb{H}^n



- Stochastic differential equations.

$$\left\{ \begin{array}{l} dX_1(t) = X_n(t)dB_1 \\ \vdots \\ dX_{n-1}(t) = X_n(t)dB_{n-1} \\ dX_n(t) = X_n(t)dB_n(t) - \frac{n-2}{2}X_n(t)dt, \end{array} \right.$$

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- System of SDE has unique strong solution and it is called hyperbolic Brownian motion.

- Solution

$$\left\{ \begin{array}{l} X_1(t) = x_1 + \int_0^t x_n \exp(B_n(u) - (n-1)/2u) dB_1(u) \\ \vdots \\ X_{n-1}(t) = x_{n-1} + \int_0^t x_n \exp(B_n(u) - (n-1)/2u) dB_{n-1}(u) \\ X_n(t) = x_n \exp(B_n(t) - (n-1)/2t). \end{array} \right.$$

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- Generator = $\frac{1}{2}\Delta$

- (1996) J.C. Gruet,
transition density function - formula

Previous results

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Poisson kernel of the ball - formula
- (2012) J. Małeckci, GS
Green function and Poisson kernel of the set
 $\{x \in \mathbb{H}^n : x_1 > 0\}$ for HBM with drift - formulas and estimates

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Notation:

- $G_U^E(x, y)$, $P_U^E(x, y)$ - Green function and Poisson kernel of a set U for standard (euclidean) Brownian motion.

We will show relationship between n - dimensional hyperbolic Brownian motion and $2n$ - dimensional standard Brownian motion

Notation:

- $G_U^E(x, y), P_U^E(x, y)$ - Green function and Poisson kernel of a set U for standard (euclidean) Brownian motion.
- $G_U^H(x, y), P_U^H(x, y)$ - Green function and Poisson kernel of a set U for hyperbolic Brownian motion (with respect to the Lebesgue measure).

Let consider sets of the form

$$\begin{aligned}\{x \in \mathbb{H}^n : (x_1, \dots, x_{n-1}) \in D\} &= D \times (0, \infty), \\ \{x \in \mathbb{R}^{2n} : (x_1, \dots, x_{n-1}) \in D\} &= D \times \mathbb{R}^{n+1},\end{aligned}$$

where $D \subseteq \mathbb{R}^{n-1}$.

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Let $x \in \mathbb{R}^m$, $m > n$. We denote

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Then for $x, y \in D \times \mathbb{R}^{n+1}$ we have

$$\begin{aligned}G_{D \times \mathbb{R}^{n+1}}^E(x, y) &= G_{D \times \mathbb{R}^{n+1}}^E(x^-, y^-, |x^+ - y^+|^2), \\ P_{D \times \mathbb{R}^{n+1}}^E(x, y) &= P_{D \times \mathbb{R}^{n+1}}^E(x^-, y^-, |x^+ - y^+|^2).\end{aligned}$$

Theorem (GS 2012)

For $x, y \in D \times (0, \infty)$ we have

$$G_{D \times (0, \infty)}^H(x, y) = \frac{2\pi^{n/2} x_n y_n^{n-1}}{\Gamma(n/2)} \int_{-1}^1 (1 - u^2)^{n/2-1} \times$$

$$G_{D \times \mathbb{R}^{n+1}}^E(x^-, y^-, x_n^2 + y_n^2 + 2x_n y_n u) du.$$

Relationships

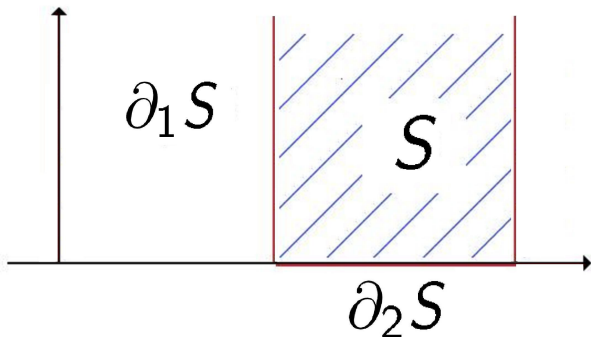
We divide the boundary ∂S of the set $S = D \times (0, \infty) \subseteq \mathbb{H}^n$, $D \subseteq \mathbb{R}^{n-1}$ into two parts

- $\partial_1 S = \partial D \times (0, \infty)$

It is a boundary S as a subset of \mathbb{H}^n .

- $\partial_2 S = D \times \{0\}$

It is not a subset of \mathbb{H}^n , but hyperbolic Brownian motion can hit it in infinite time.



Theorem (GS 2012)

For $x \in S, y \in \partial_1 S$ we have

$$P_S^H(x, y) = \frac{2\pi^{n/2} y_n x_n^{n-1}}{\Gamma(n/2)} \int_{-1}^1 (1 - u^2)^{n/2-1} \times$$

$$P_{D \times \mathbb{R}^{n+1}}^E(x^-, y^-, x_n^2 + y_n^2 + 2x_n y_n u) du.$$

For $x \in S, y \in \partial_2 S$ we have similar formula. However, it is more complicated and less useful.

For $a > 0$ we define

$$S_a = \{x \in \mathbb{H}^n : x_1 \in (0, a)\}.$$

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We divide the boundary as before:

$$\begin{aligned}\partial_1 S_a &= \{0, a\} \times \mathbb{R}^{n-2} \times (0, \infty), \\ \partial_2 S_a &= (0, a) \times \mathbb{R}^{n-2} \times \{0\}.\end{aligned}$$

Theorem (GS 2012)

For $x \in S_a$, $y \in \partial_1 S_a$ we have

$$P_{S_a}^H(x, y) = \frac{(-2)^{1-n} \pi^{\frac{n}{2}}}{a \Gamma(\frac{n}{2})} \frac{x_n}{y_n^{n-1}} \int_{-1}^1 (1-u^2)^{\frac{n-2}{2}} \frac{\partial^{n-1}}{\partial u^{n-1}} \frac{\sin\left(\frac{\pi(x_1-y_1)}{a}\right) du}{\cosh\left(\frac{\pi}{a} \sqrt{2x_n y_n (\cosh \tilde{\rho} + u)}\right) - \cos\left(\frac{\pi}{a}(x_1 - y_1)\right)},$$

Theorem (GS 2012)

and for $y \in \partial_2 S_a$ we have

$$P_{S_a}^H(x, y) = \frac{2x_n^{n-1}}{\Gamma\left(\frac{n+1}{2}\right)} \frac{(-1)^n}{\pi^{\frac{n-1}{2}}} \times \\ \frac{\partial^{n-1}}{\partial \xi^{n-1}} \ln \left[1 - \frac{2 \sin x_1 \sin y_1}{\cosh\left(\frac{\pi}{a} \sqrt{\xi}\right) + \cos\left(\frac{\pi(x_1 - y_1)}{a}\right)} \right]_{\xi = |\tilde{x} - \tilde{y}|^2},$$

where $\tilde{x} = (x_2, x_3, \dots, x_n)$.

Estimates in euclidean case

$$P_{(0,a) \times \mathbb{R}^{2n-1}}^E(x, y) \asymp \frac{x_1(a - x_1) \exp\left(-\frac{\pi}{a} |x - y|\right) a^{n+1} + |x - y|^{n+1}}{a^{n+2} |x - y|^{2n}}.$$

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Theorem (GS 2012)

Let $x \in (0, a) \times \mathbb{H}^{n-1}$. For $y \in \partial_1 S_a$ we have

$$P_{S_a}^H(x, y) \underset{c}{\asymp} \frac{x_1(a-x_1) e^{-\frac{a}{\pi}|x-y|}}{a^2 |x-y|^n} \left(\frac{x_n}{y_n} \right)^{n/2} \frac{a^{n+1} + |x-y|^{n+1}}{(a^2 \cosh \rho + |x-y|)^{n/2}},$$

Theorem (GS 2012)

and for $y \in \partial_2 S_a$ we have

$$P_{S_a}^H(x, y) \asymp \frac{c x_n^{n-1}}{a^{n+3}} \exp\left(-\frac{\pi}{a} |x - y|\right) \frac{x_1 y_1 (a - x_1)(a - y_1)}{|\tilde{x} - \tilde{y}|^2 + (a - |x_1 + y_1 - a|)^2} \frac{a^{n+1} + |x - y|^{n+1}}{|x - y|^{2n-2}}.$$

The constant c depends only on dimension n .

We define the functional

$$A_{x_n}(t) = x_n^2 \int_0^t \exp(2X_n(u)) du.$$

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We have

$$X(t) \stackrel{(d)}{=} Y(A_{x_n}(t)),$$

where

$$Y(t) = \left(\beta_1(t), \dots, \beta_{n-1}(t), R^{(-\frac{n-1}{2})}(t) \right).$$

Here process $R^{(-\frac{n-1}{2})}(t)$ is the Bessel process starting from $x_n > 0$ and independent of the process $\beta(t)$. The integral functional $A_{x_n}(t)$ is independent of $\beta(t)$ too.

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Remark Poisson kernels for processes $X(t)$ and $Y(t)$ are the same.

Process $Y(t)$ without last coordinate is just standard $n - 1$ - dimensional Brownian motion. For the last coordinate we have

$$\mathbb{P}^x \left(R_t^{(-\frac{n-1}{2})} \in dy \right) = x^{n-1} \mathbb{P}^x \left(\left(R_t^{(\frac{n-1}{2})} \right)^{-n-1}, R_t^{(\frac{n-1}{2})} \in dy \right).$$

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$R_t^{(\frac{n-1}{2})}$ is the absolute value of $n + 1$ - dimensional Brownian motion.