Dirichlet Heat Kernel Estimates for Relativistic Stable Processes

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References

This talk is based on the following joint papers with Zhen-Qing Chen and Panki Kim.

- CKS1 Sharp Heat Kernel Estimates for Relativistic Stable Processes in Open Sets. Ann. Probab. 40 (1) (2012), 213–244.
- CKS2 Global heat kernel estimates for relativistic stable processes in half-space-like open sets. *Potential Anal.*, **36** (2012) 235–261.
- CKS3 Global heat kernel estimate for relativistic stable processes in exterior open sets. *J. Funct. Anal.*, **263** (2012), 448–475.
- **CKS4** Dirichlet heat kernel estimates for rotationally symmetric Lévy processes. In preparation.
- **CKS5** Dirichlet heat kernel estimates for subordinate Brownian motions with Gaussian components. In preparation.

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- 2 Relativistic stable processes
- Estimates in General smooth open sets
- Dirichlet heat kernel estimates for subordinate Brownian motions

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- 5 Dirichlet Heat kernel estimates in half-space-like open set
- Dirichlet Heat kernel estimates in exterior open set

Suppose that *X* is a symmetric Markov process on (a subset of) \mathbb{R}^d with transition density p(t, x, y) and generator \mathcal{L} . p(t, x, y) is also the fundamental solution of $\partial_t u = \mathcal{L}u$ and so it is also called the heat kernel of \mathcal{L} . In general, there is no explicit formula for p(t, x, y). Thus establishing sharp two-sided estimates for p(t, x, y) is a fundamental problem.

Two-sided heat kernel estimates for diffusions in \mathbb{R}^d have a long history and many beautiful results have been established. Among the main contributors are: D. G. Aronson, J. Nash, E. B. Davies.

Due to the complication near the boundary, two-sided estimates on the transition density of killed diffusions in a domain D (equivalently, the Dirichlet heat kernel) have been established only recently. See Davies (87), Zhang (02) for the case of bounded $C^{1,1}$ domains.

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The infinitesimal generator of a discontinuous Markov process in \mathbb{R}^d is no longer a differential operator but rather a non-local (or integro-differential) operator \mathcal{L} . For instance, the infinitesimal generator of a rotationally symmetric α -stable process in \mathbb{R}^d with $\alpha \in (0, 2)$ is a fractional Laplacian operator $\Delta^{\alpha/2} := -(-\Delta)^{\alpha/2}$.

Recently in [CKS10, JEMS], we obtained sharp two-sided estimates for the heat kernel of the fractional Laplacian $\Delta^{\alpha/2}$ in *D* with zero exterior condition (or equivalently, the transition density function of the symmetric α -stable process killed upon exiting *D*) for any $C^{1,1}$ open set $D \subset \mathbb{R}^d$ with $d \ge 1$. As far as we know, this was the first time sharp two-sided estimates were established for Dirichlet heat kernels of non-local operators. The infinitesimal generator of a discontinuous Markov process in \mathbb{R}^d is no longer a differential operator but rather a non-local (or integro-differential) operator \mathcal{L} . For instance, the infinitesimal generator of a rotationally symmetric α -stable process in \mathbb{R}^d with $\alpha \in (0, 2)$ is a fractional Laplacian operator $\Delta^{\alpha/2} := -(-\Delta)^{\alpha/2}$.

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In this talk, I will present sharp two-sided estimates on the Dirichlet heat kernels of relativistic stable processes in $C^{1,1}$ domains of \mathbb{R}^d .

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- **6** Dirichlet Heat kernel estimates in exterior open set

For $\alpha \in (0, 2]$, a symmetric α -stable process X on \mathbb{R}^d is a Lévy process such that for any $t \ge 0$ and $\xi \in \mathbb{R}^d$

$$\mathbb{E}\left[\exp\left(i\xi\cdot(X_t-X_0)\right)\right]=\exp\left(-t|\xi|^{\alpha}\right).$$

When $\alpha = 2$, it reduces to a Brownian motion.

The infinitesimal generator of a symmetric α -stable process Y in \mathbb{R}^d is the fractional Laplacian $\Delta^{\alpha/2}$, which can be written as

$$\Delta^{\alpha/2} u(x) = \lim_{\varepsilon \downarrow 0} \int_{\{y \in \mathbb{R}^d : |y-x| > \varepsilon\}} (u(y) - u(x)) \frac{\mathcal{A}(d, \alpha)}{|x-y|^{d+\alpha}} \, dy,$$

where $\mathcal{A}(d, \alpha) := \alpha 2^{\alpha-1} \pi^{-d/2} \Gamma(\frac{d+\alpha}{2}) \Gamma(1-\frac{\alpha}{2})^{-1}$.

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Let p(t, x, y) be the transition density of X. When $d > \alpha$, the potential density (also called the Green function) of X

$$G(x,y) = \int_0^\infty p(t,x,y) dt = C(d,\alpha) \frac{1}{|x-y|^{d-\alpha}}$$

which is the Riesz kernel.

Symmetric stable processes have some nice properties. For example it satisfies the following scaling property: For any a > 0, $\{a^{-1/\alpha}(X_{at} - X_0) : t \ge 0\}$ has the same law as $\{X_t - X_0 : t \ge 0\}$. In terms of the transition density, this means

$$p(t, x, y) = a^d p(at, a^{1/\alpha}x, a^{1/\alpha}y).$$

However, a symmetric α -stable process, for $\alpha \in (0, 2)$, always have infinite variance. When $\alpha \in (0, 1]$, it also have infinite mean.

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For any $m \ge 0$, a relativistic α -stable process X^m on \mathbb{R}^d with weight m is a Lévy process such that for any $t \ge 0$ and $\xi \in \mathbb{R}^d$

$$\mathbb{E}\left[\exp\left(i\xi\cdot(X_t^m-X_0^m)\right)\right]=\exp\left(-t\left(\left(|\xi|^2+m^{2/\alpha}\right)^{\alpha/2}-m\right)\right).$$

When m = 0, X^0 is simply a (rotationally) symmetric α -stable process on \mathbb{R}^d . The infinitesimal generator of X^m is

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When $\alpha = 1$, the infinitesimal generator reduces to

$$m-\sqrt{-\Delta+m^2}.$$

This operator was used by E. Lieb and his followers in studying the stability of matter.

Let $p^m(t, x, y)$ be the transition density of X^m . The the function, called the 1-potential density of X^1 :

$$\int_0^\infty e^{-t} p^1(t, x, y) dt$$

is the Bessel kernel.

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The Lévy measure of X^m has a density

$$J^m(\mathbf{x}, \mathbf{y}) = \mathcal{A}(\mathbf{d}, \alpha) |\mathbf{x} - \mathbf{y}|^{-\mathbf{d} - \alpha} \psi(\mathbf{m}^{1/\alpha} |\mathbf{x} - \mathbf{y}|)$$

where

$$\psi(r) := \int_0^\infty s^{\frac{d+\alpha}{2}-1} e^{-\frac{s}{4}-\frac{r^2}{s}} \, ds,$$

which is decreasing and is a smooth function of r^2 satisfying $\psi(r) \leq 1$ and

$$\psi(r) \asymp \phi(r) := e^{-r} (1 + r^{(d+\alpha-1)/2})$$
 on $[0, \infty)$.

For m > 0, X^m has moments of all orders, and it even has some exponential moments. In a small scale, X^m behaves like X^0 , while in a larger scale, X^m behaves like Brownian motion.

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For m > 0, X^m does not satisfy any scaling invariance property. However, it does satisfy some sort of approximate scaling property.

Two-sided estimates on p(t, x, y) is classical. But two-sided estimates on $p^m(t, x, y)$ is more recent.

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For any
$$m, c > 0$$
, we define a function $\widetilde{\Psi}_{d,\alpha,m,c}(t, x, y)$ on
 $(0, \infty) \times \mathbb{R}^d \times \mathbb{R}^d$ by
 $\widetilde{\Psi}_{d,\alpha,m,c}(t, x, y)$

$$= \begin{cases} t^{-d/\alpha} \wedge tJ^m(x, y), & \forall t \in (0, 1/m]; \\ m^{d/\alpha - d/2}t^{-d/2}\exp\left(-c^{-1}(m^{1/\alpha}|x-y| \wedge m^{2/\alpha - 1}\frac{|x-y|^2}{t})\right), & \forall t \in (1/m, \infty). \end{cases}$$

Theorem [Chen-Kim-Kumagai], [CKS1

 $c_1^{-1}\widetilde{\Psi}_{d,\alpha,m,1/C_1}(t,x,y) \le p^m(t,x,y) \le c_1\widetilde{\Psi}_{d,\alpha,m,C_1}(t,x,y).$

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Relativistic stable process in *D*

For any open set *D*, we use $\tau_D^m := \inf\{t > 0 : X_t^m \notin D\}$ to denote the first exit time from *D* by X^m , and $X^{m,D}$ to denote the subprocess of X^m killed upon exiting *D* (or, the killed relativistic stable process in *D* with mass *m*). We will use $p_D^m(t, x, y)$ to denote the transition density of $X^{m,D}$.

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Theorem [CKS1]

Suppose that D is a $C^{1,1}$ open set. (i) For any $m \in (0, M]$ and $(t, x, y) \in (0, T] \times D \times D$,

$$\frac{1}{C_1} \left(1 \wedge \frac{\delta_D(\mathbf{x})^{\alpha/2}}{\sqrt{t}} \right) \left(1 \wedge \frac{\delta_D(\mathbf{y})^{\alpha/2}}{\sqrt{t}} \right) \left(t^{-d/\alpha} \wedge \frac{t\phi(m^{1/\alpha}|\mathbf{x} - \mathbf{y}|)}{|\mathbf{x} - \mathbf{y}|^{d+\alpha}} \right)$$

$$\leq p_D^m(t, \mathbf{x}, \mathbf{y}) \leq$$

$$C_1 \left(1 \wedge \frac{\delta_D(\mathbf{x})^{\alpha/2}}{\sqrt{t}} \right) \left(1 \wedge \frac{\delta_D(\mathbf{y})^{\alpha/2}}{\sqrt{t}} \right) \left(t^{-d/\alpha} \wedge \frac{t\phi(m^{1/\alpha}|\mathbf{x} - \mathbf{y}|/(16))}{|\mathbf{x} - \mathbf{y}|^{d+\alpha}} \right)$$

where $\phi(r) = e^{-r}(1 + r^{(d+\alpha-1)/2})$. (ii) Suppose in addition that *D* is bounded. for any $m \in (0, M]$ and $(t, x, y) \in [T, \infty) \times D \times D$,

$$\mathcal{P}_D^m(t, \mathbf{x}, \mathbf{y}) \asymp \mathbf{e}^{-t \, \lambda_1^{\alpha, m, D}} \, \delta_D(\mathbf{x})^{\alpha/2} \, \delta_D(\mathbf{y})^{\alpha/2},$$

where $\lambda_1^{\alpha,m,D} > 0$ is the smallest eigenvalue of the restriction of $(m^{2/\alpha} - \Delta)^{\alpha/2} - m$ in *D* with zero exterior condition.

Our estimates are uniform in *m* in the sense that the constants are independent of $m \in (0, M]$. Letting $m \downarrow 0$ recovers the below sharp heat kernel estimates for symmetric α -stable processes obtained in [CKS, JEMS10].

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Difficulties and Ingredients

- two-sided estimates on p^m.
- the approximate scaling property
- the Lévy density of X^m does not have a simple form and has exponential decay rate as oppose to the polynomial decay rate of the Lévy density of symmetric stable process
- uniform Boundary Harnack principle and parabolic Harnack principle
- There exist positive constants R₀ and C > 1 depending only on d and α such that for any m ∈ (0,∞), any ball B of radius r ≤ R₀m^{-1/α},

$$C^{-1}G_B(x,y) \leq G_B^m(x,y) \leq CG_B(x,y), \quad x,y \in B.$$

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Stable processes and relativistic stable processes are all examples of subordinate Brownian motions.

Suppose that
$$S = \{S_t : t \ge 0\}$$
 is a subordinator with Laplace exponent ϕ :
 $\mathbb{P} e^{-\lambda S_t} = e^{-t\phi(\lambda)}$

Suppose that $B = \{B_t : t \ge 0\}$ is *d*-dimensional Brownian motion independent of the subordinator *S*. Then the process $X = \{X_t : t \ge 0\}$ defined by $X_t = B_{S_t}$ is called a subordinate Brownian motion and it is a rotationally symmetric Lévy process.

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In this section we will always assume that ϕ is a complete Bernstein function, that is, the Lévy measure μ has a density $\mu(t)$ which is completely monotone $((-1)^n \mu^{(n)}(t) \ge 0$ for n = 1, 2, ...). *X* has a transition density p(t, x, y) with respect to the Lebesgue measure.

For any open set $D \subset \mathbb{R}^d$, we will use X^D to denote the process obtained from X by killing it upon exiting from D. The process X^D has a transition density $p^D(t, x, y)$ with respect to the Lebesgue measure on D. The main results of this section are sharp two-sided estimates on $p^D(t, x, y)$.

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For the first result in this section, we will assume the following **(H)**: There exist constants $\delta_1, \delta_2 \in (0, 1)$, $a_1, a_2 > 0$ and $R_0 > 0$ such that

(LSC)	$\phi(\lambda r) \geq a_1 \lambda^{\delta_1} \phi(r),$	$\lambda \ge 1, r \ge 1/R_0^2$
(USC)	$\phi(\lambda r) \leq a_2 \lambda^{\delta_2} \phi(r),$	$\lambda \geq 1, r \geq 1/R_0^2.$

Note that it follows from (USC) that ϕ has no drift.

Definition

Suppose R > 0 and $\kappa \in (0, 1)$. An open set $D \subset \mathbb{R}^d$ is called κ -fat if there is R > 0 such that for every $x \in \overline{D}$ and all $r \in (0, R]$, $D \cap B(x, r)$ contains a ball of radius κr . The pair (R, κ) is called the characteristics of the κ -fat open set D.

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Theorem

Suppose that *D* is a bounded κ -fat open set in \mathbb{R}^d .

For every *T* > 0, there exist *c_i* = *c_i*(*R*, κ, *T*, *d*, φ) > 1, *i* = 1, 2 such that for 0 < *t* ≤ *T*, *x*, *y* ∈ ℝ^d,

$$p_{\mathsf{D}}(t,x,y) \leq c_1 \mathbb{P}_x(au_{\mathsf{D}} > t) \mathbb{P}_y(au_{\mathsf{D}} > t) p(t,c_2^{-1}x,c_2^{-1}y)$$

and

$$c_1^{-1}\mathbb{P}_x(au_D > t)\mathbb{P}_y(au_D > t) p(t, c_2 x, c_2 y) \leq p_D(t, x, y)$$

For every T > 0, there is a constant c₃ ≥ 1 depending only on diam(D), T, R, κ, d and φ so that for all (t, x, y) ∈ [T, ∞) × D × D,

$$\begin{array}{lll} \rho_D(t,x,y) &\geq & c_3^{-1}\,\mathbb{P}_x(\tau_D>1)\mathbb{P}_y(\tau_D>1)e^{-t\lambda_1}\\ \rho_D(t,x,y) &\leq & c_3\,\mathbb{P}_x(\tau_D>1)\mathbb{P}_y(\tau_D>1)e^{-t\lambda_1}, \end{array}$$

where $-\lambda_1 < 0$ is the largest eigenvalue of the generator of X^D .

We can actually prove the first part of the above theorem without the boundedness assumption on *D*, but with a little extra condition on ϕ which we think is not necessary.

In the above theorem, we do not really need *X* to be a subordinate Brownian motion. What we really need is that *X* is a rotationally symmetric, purely discontinuous Lévy process whose Lévy density is comparable to that of a subordinate Brownian motion satisfying the assumptions of the above theorem.

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Theorem (C^{1.1} case)

Suppose that

$$\phi(\lambda) symp \lambda^{lpha/2} \ell(\lambda), \qquad \lambda o \infty,$$

where $\alpha \in (0, 2)$ and ℓ is a positive function which is slowing varying at infinity. Let *D* be a bounded $C^{1,1}$ open subset of \mathbb{R}^d with characteristics (R_0, Λ_0) . (1) For every T > 0, there exist $c_j = c_j(R_0, \Lambda_0, T, d, \alpha, \ell) \ge 1, j = 1, 2$, such that for all $(t, x, y) \in (0, T] \times D \times D$,

$$\begin{split} c_{1}^{-1} \left(1 \wedge \frac{1}{t\phi(\delta_{D}^{-2}(x))} \right)^{1/2} \left(1 \wedge \frac{1}{t\phi(\delta_{D}^{-2}(y))} \right)^{1/2} \left(\frac{1}{(\Phi^{-1}(t))^{d}} \wedge tJ(c_{2}x, c_{2}y) \right) \\ &\leq \rho_{D}(t, x, y) \leq \\ c_{1} \left(1 \wedge \frac{1}{t\phi(\delta_{D}^{-2}(x))} \right)^{1/2} \left(1 \wedge \frac{1}{t\phi(\delta_{D}^{-2}(y))} \right)^{1/2} \left(\frac{1}{(\Phi^{-1}(t))^{d}} \wedge tJ(x/C_{2}, y/C_{2}) \right) \\ &\text{where } \Phi(r) = \frac{1}{\phi(r^{-2})}. \end{split}$$

Theorem (Cont)

(2) For every T > 0, there is a constant $c_3 \ge 1$ depending only on diam(D), R_0 , Λ_0 , d, α , ℓ and T so that for all $(t, x, y) \in [T, \infty) \times D \times D$,

$$egin{array}{rcl} p_D(t,x,y) &\geq & c_3^{-1} \, e^{-\lambda_1 t} rac{1}{\sqrt{\phi(\delta_D^{-2}(x))}} rac{1}{\sqrt{\phi(\delta_D^{-2}(y))}} \ p_D(t,x,y) &\leq & c_3 \, e^{-\lambda_1 t} \, rac{1}{\sqrt{\phi(\delta_D^{-2}(x))}} rac{1}{\sqrt{\phi(\delta_D^{-2}(y))}}, \end{array}$$

where $-\lambda_1 < 0$ is the largest eigenvalue of the generator of X^D .

Again, we get get rid of the boundedness assumption on *D* if we can assume a little extra condition ϕ which we think is not necessary. We can also deal with the case when the subordinate Brownian motion *X* has a Gaussian component when *D* is a $C^{1,1}$ open set.

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Dirichlet Heat kernel estimates in half-space-like open set

Dirichlet Heat kernel estimates in exterior open set

A half-space is any set which, after isometry, can be written as $\{(x_1, \ldots, x_d) : x_d > 0\}.$

An open set *D* is said to be half-space-like if, after isometry, $H_a \subset D \subset H_b$ for some real numbers a > b. Here for any real number $a, H_a := \{(x_1, \dots, x_d) : x_d > a\}$. H_0 will be simply written as *H*.

For any
$$m, c > 0$$
, define

$$\begin{aligned}
\Psi_{d,\alpha,m,c}(t, x, y) &= \begin{cases} t^{-d/\alpha} \wedge \frac{t\phi(c^{-1}m^{1/\alpha}|x-y|)}{|x-y|^{d+\alpha}} & t \in (0, 1/m], \\ m^{d/\alpha - d/2}t^{-d/2}\exp\left(-c^{-1}(m^{1/\alpha}|x-y| \wedge m^{2/\alpha - 1}\frac{|x-y|^2}{t})\right) & t \in (1/m, \infty), \end{cases}
\end{aligned}$$
where $\phi(r) = e^{-r} \left(1 + r^{(d+\alpha - 1)/2}\right).$

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$$\begin{aligned} & \text{For any } m, c > 0, \text{ define} \\ & \Psi_{d,\alpha,m,c}(t, x, y) \\ & := \begin{cases} t^{-d/\alpha} \wedge \frac{t\phi(c^{-1}m^{1/\alpha}|x-y|)}{|x-y|^{d+\alpha}} & t \in (0, 1/m], \\ m^{d/\alpha - d/2}t^{-d/2}\exp\left(-c^{-1}(m^{1/\alpha}|x-y| \wedge m^{2/\alpha - 1}\frac{|x-y|^2}{t})\right) & t \in (1/m, \infty), \end{cases} \\ & \text{ where } \phi(r) = e^{-r} \left(1 + r^{(d+\alpha - 1)/2}\right). \end{aligned}$$

Theorem [CKS2]

Suppose *D* is a half-space-like $C^{1,1}$ open set. For any M > 0, there exist $C_i > 1 \ge 1$, i = 1, 2, such that for all $m \in (0, M]$, (i) if $t \in (0, 1/m]$

$$\begin{split} C_1^{-1} \left(\frac{\delta_D(\boldsymbol{x})^{\alpha/2}}{\sqrt{t}} \wedge 1 \right) \left(\frac{\delta_D(\boldsymbol{y})^{\alpha/2}}{\sqrt{t}} \wedge 1 \right) \Psi_{d,\alpha,m,C_2}(t,\boldsymbol{x},\boldsymbol{y}) &\leq p_D^m(t,\boldsymbol{x},\boldsymbol{y}) \\ &\leq C_1 \left(\frac{\delta_D(\boldsymbol{x})^{\alpha/2}}{\sqrt{t}} \wedge 1 \right) \left(\frac{\delta_D(\boldsymbol{y})^{\alpha/2}}{\sqrt{t}} \wedge 1 \right) \Psi_{d,\alpha,m,1/C_2}(t,\boldsymbol{x},\boldsymbol{y}) \end{split}$$

(ii) if *t* > 1/*m*

$$\begin{split} C_1^{-1} \left(\frac{m^{(2-\alpha)/2\alpha} \delta_D(\boldsymbol{x}) + \delta_D(\boldsymbol{x})^{\alpha/2}}{\sqrt{t}} \wedge 1 \right) \left(\frac{m^{(2-\alpha)/2\alpha} \delta_D(\boldsymbol{y}) + \delta_D(\boldsymbol{y})^{\alpha/2}}{\sqrt{t}} \wedge 1 \right) \\ & \times \Psi_{d,\alpha,m,C_2}(t,\boldsymbol{x},\boldsymbol{y}) \\ & \leq p_D^m(t,\boldsymbol{x},\boldsymbol{y}) \leq \\ C_1 \left(\frac{m^{(2-\alpha)/2\alpha} \delta_D(\boldsymbol{x}) + \delta_D(\boldsymbol{x})^{\alpha/2}}{\sqrt{t}} \wedge 1 \right) \left(\frac{m^{(2-\alpha)/2\alpha} \delta_D(\boldsymbol{y}) + \delta_D(\boldsymbol{y})^{\alpha/2}}{\sqrt{t}} \wedge 1 \right) \\ & \times \Psi_{d,\alpha,m,1/C_2}(t,\boldsymbol{x},\boldsymbol{y}). \end{split}$$

These estimates are new even when *D* is the upper half space *H*. Observe that although *H* is invariant under scaling, global two-sided estimates on $p_H^m(t, x, y)$ can not be derived through a scaling argument from the short time estimates which hold only for $m \in (0, M]$ and $t \in (0, T]$.

For a fixed half-space-like $C^{1,1}$ open set D with $C^{1,1}$ characteristics (R, Λ_0) and $H_a \subset D \subset H_b$, mD is still a half-space-like $C^{1,1}$ open set but with $C^{1,1}$ -characteristics $(mR, \Lambda_0/m)$ and $H_{ma} \subset mD \subset H_{mb}$. So we can not use the scaling property

$$p_D^m(t, x, y) = m^{d/\alpha} p_{m^{1/\alpha}D}^1(mt, m^{1/\alpha}x, m^{1/\alpha}y)$$

to obtain sharp two-sided estimates for $p_D^m(t, x, y)$ that are uniform in $m \in (0, M]$ from that of $p_D^1(t, x, y)$.

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A major part is to derive global sharp two-sided heat kernel estimates for X^m in a half-space.

Then, we use the push-inward technique developed in [Chen-Tokle, PTRF 2011] to extend it to half-space-like $C^{1,1}$ open sets.

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5 Dirichlet Heat kernel estimates in half-space-like open set

Dirichlet Heat kernel estimates in exterior open set

An open set *D* in \mathbb{R}^d is called an exterior open set if D^c is compact.

Theorem [CKS3]

Suppose that $d \ge 3$, M > 0 and D is an exterior $C^{1,1}$ open set in \mathbb{R}^d . Then there are constants $c_i > 1$, i = 1, 2, such that for every $m \in (0, M]$, t > 0 and $(x, y) \in D \times D$,

$$p_D^m(t, x, y) \le c_1 \left(1 \land \frac{\delta_D(x)}{1 \land t^{1/\alpha}} \right)^{\alpha/2} \left(1 \land \frac{\delta_D(y)}{1 \land t^{1/\alpha}} \right)^{\alpha/2} \Psi_{d, \alpha, m, c_2}(t, x, y)$$

and

$$p_D^m(t,x,y) \ge c_1^{-1} \left(1 \wedge \frac{\delta_D(x)}{1 \wedge t^{1/\alpha}}\right)^{\alpha/2} \left(1 \wedge \frac{\delta_D(y)}{1 \wedge t^{1/\alpha}}\right)^{\alpha/2} \Psi_{d,\alpha,m,1/o_2}(t,x,y).$$

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and

$$p_D^m(t, x, y) \ge c_1^{-1} \left(1 \wedge \frac{\delta_D(x)}{1 \wedge t^{1/\alpha}}\right)^{\alpha/2} \left(1 \wedge \frac{\delta_D(y)}{1 \wedge t^{1/\alpha}}\right)^{\alpha/2} \Psi_{d,\alpha,m,1/c_2}(t, x, y).$$

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The reason that we assume $d \ge 3$ is that we used the transience of X^m . By Chung-Fuch's criterion for Lévy processes, X^m is transient if and only if $d \ge 3$.

The large time upper bound is relatively easy to establish. The main difficulty is in establishing the large time lower bound.

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Thank you!

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